# THE FORCE CONSTANTS OF PHOSPHOROUS OXYCHLORIDE 

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Received October 19, 1955

## Summary

The force constants of phosphorous oxychloride molecule, $\mathrm{OPCl}_{3}$, have been calculated using the matrix method of Wilson for a valence force potential function with three interaction terms. The values of the force constants are discussed in relation to those of related molecules.

## Introduction

Molecules of the type $\mathrm{XYZ}_{3}$ which belong to the point group $\mathrm{C}_{3 \mathrm{~g}}$ have been studied extensively regarding their vibrational spectra and molecular structure. The oxyhalides of phosphorus form an interesting series in this group. The molecular structure of phosphorous oxychloride with which this paper is concerned has been determined from electron diffraction measurements by Brockway and Beach (1938) and more recently from microwave spectrum by Quitman et al. (1952). The vibrational spectrum has been studied by Langseth (1931), Cabannes and Rousset (1933), Thatte (1936), Simon and Schulze (1937) and Gerding and Westrick (1942). As here appears to be no calculation of the force constants of this molecule it is considered worth while to evaluate the same by Wilson's matrix method using the molecular dimensions of Quitman et al.

## Symmetry Co-ordinates

The five atomic molecule of point group $\mathrm{C}_{30}$ will have nine fundamental vibration frequencies three of which belong to the totally symmetric non-degenerate Class $\mathrm{A}_{1}$, and three to the doubly degenerate Class E. The symmetry co-ordinates for these two species can he formulated in terms of the internal co-ordinates as follows:-

For Type A-

$$
\begin{aligned}
& \mathrm{R}_{1}=\Delta \mathrm{D} \\
& \mathrm{R}_{2}=\frac{1}{\sqrt{3}}\left(\Delta d_{1}+\Delta d_{2}+\Delta d_{3}\right) \\
& \mathrm{R}_{3}=\frac{1}{\sqrt{6}}\left(\Delta a_{12}+\Delta a_{23}+\Delta a_{31}-\Delta \beta_{1}-\Delta \beta_{2}-\Delta \beta_{3}\right)
\end{aligned}
$$

For Type E-

$$
\begin{aligned}
& \mathrm{R}_{1 a}=\frac{1}{\sqrt{6}}\left(2 \Delta d_{1}-\Delta d_{2}-\Delta d_{3}\right) \\
& \mathrm{R}_{1 b}=\frac{1}{\sqrt{2}}\left(\Delta d_{2}-\Delta d_{3}\right) \\
& \mathrm{R}_{2 a}=\frac{1}{\sqrt{6}}\left(2 \Delta \beta_{1}-\Delta \beta_{9}-\Delta \beta_{3}\right) \\
& \mathrm{R}_{2 b}=\frac{1}{\sqrt{2}}\left(\Delta \beta_{2}-\Delta \beta_{3}\right) \\
& \mathrm{R}_{s a}=\frac{1}{\sqrt{6}}\left(2 \Delta a_{23}-\Delta a_{12}-\Delta a_{31}\right) \quad \mathrm{R}_{3 b}=\frac{1}{\sqrt{2}}\left(\Delta \alpha_{31}-\Delta \alpha_{12}\right)
\end{aligned}
$$

D is the O-P distance, $d_{i}$ the $\mathrm{P}-\mathrm{Cl}_{i}$ distance, $a_{i j}$ the angle $\mathrm{Cl}_{i} \mathrm{PCl}$, and $\beta_{i}$ the angle $\mathrm{Cl}_{i} \mathrm{PO}$ the values being taken at the equilibrium position of the molecule.

## The Kinetic Energy Matrix

The G matrix associated with the kinetic energy of the molecule was constructed by first obtaining the $\vec{s}$ vectors in terms of the unit vectors directed along the bonds, the equilibrium bond lengths and bond angles in the usual manner and making use of the relations

$$
\mathrm{G}_{j l}=\sum_{p} \mu_{p} g_{n} \overrightarrow{\mathrm{~S}}_{j}^{t} \cdot \overrightarrow{\mathrm{~S}}_{l}^{t}
$$

for non-degenerate vibrations and

$$
\mathrm{G}_{j t}=\frac{1}{2} \sum_{p} \mu_{p} g_{p}\left(\overrightarrow{\mathrm{~S}}_{j a}^{t} \cdot \overrightarrow{\mathrm{~S}}_{i a}{ }^{t}+\overrightarrow{\mathrm{S}}_{j b}{ }^{t} \cdot \overrightarrow{\mathrm{~S}}_{i b}{ }^{t}\right)
$$

for doubly degenerate vibrations where

$$
\overrightarrow{\mathbf{S}}_{j}^{t} \cdot=\sum_{k} \mathrm{U}_{j k} \cdot \overrightarrow{\mathbf{s}}_{k t}
$$

$\mathrm{U}_{j k}$ being the elements of the matrix connecting the internal co-ordinates and the symmetry co-ordinates [Meister and Cleveland (1946)]. With the molecular dimensions $\mathrm{D}=1.45 \mathrm{~A}, d=1.99 \mathrm{~A}, \alpha=104^{\circ}$ and $B=115^{\circ}$. The G matrices are obtained as given below.

For Type A-

$$
\begin{array}{rrr}
5.705 \times 10^{22} & -1.42 \times 10^{22} & 2.394 \times 10^{30} \\
2.569 \times 10^{22} & -1.966 \times 10^{30} \\
& 4.1796 \times 10^{38}
\end{array}
$$

For Type E-

| $4.184 \times 10^{22}$ | $-2.453 \times 10^{30}$ | $1.625 \times 10^{30}$ |
| :---: | :---: | :---: |
|  | $5.627 \times 10^{38}$ | $\cdots 1.354 \times 10^{38}$ |
|  |  | $2.058 \times 10^{38}$ |

## The Potential Energi Matrix

As there are only six frequencies only six force constants can be evaluated unless some assumptions are made in order to include more force constants. In the case of the molecule phosphorous oxyfluoride Goutousky and Liehr (1952) have assumed a valence potential function with three interaction terms. A similar potential function is assumed in this case also. It is given by

$$
\begin{aligned}
2 \mathrm{~V}= & f_{\mathrm{n}} \Delta \mathrm{D}^{2}+f_{d}\left(\Delta d_{1}^{2}+\Delta d_{2}^{2}+d_{3}^{2}\right)+2 f_{d a}\left(\Delta d_{1} \Delta d_{2}+\Delta d_{2} \dot{\Delta} d_{3}\right. \\
& \left.+\Delta d_{3} \Delta d_{1}\right)+d^{2} f_{a}\left(\Delta a_{12}{ }^{2}+\Delta a_{23}{ }^{2}+\Delta \dot{a}_{31}{ }^{2}\right)+d^{2} ر_{\beta}\left(\Delta \beta_{1}^{2}+\Delta \beta_{2}^{2}+\Delta \beta_{3}{ }^{2}\right) \\
& +2 d^{2} f_{a \alpha}\left(\Delta a_{12} \Delta a_{23}+\Delta a_{23} \Delta a_{31}+\Delta \alpha_{31} \Delta a_{12}\right)+2 d^{2} f_{\beta \beta}\left(\Delta \beta_{1} \Delta \beta_{2}\right. \\
& \left.+\Delta \beta_{2} \Delta \beta_{3}+\Delta \beta_{3} \Delta \beta_{1}\right)
\end{aligned}
$$

This potential function contains seven force constants and hence the angle angle interaction terms could be obtained only in combination. The F matrices for the species $A_{1}$, and $E$ are respectively

$$
\left|\begin{array}{ccc}
f_{\mathrm{o}} & f_{d}+2 f_{d d} & \ddots \\
& & \therefore \frac{1}{2} d^{2}\left(f_{\alpha}+f \beta+2 f_{a \alpha}+2 f_{\beta \beta}\right)
\end{array}\right| \ldots
$$

and

$$
f_{d}-f_{d d}: \begin{array}{cc}
d^{2}\left(f_{\beta}-f_{\beta \beta}\right) & 0  \tag{0}\\
\hdashline & d^{2}\left(f_{a}-f_{a \alpha}\right)
\end{array}
$$

From the F and G matrices the expanded form of the secular equation is obtained and the force constants were determined which provided a best fit for the observed fundamental frequencies. The observed frequencies taken from Gerding and Westrik are 1290,486 and $267 \mathrm{~cm} .^{-1}$ for Type $\mathrm{A}_{1}$, and 581,337 and $193 \mathrm{~cm} .^{-1}$ for Type E . The force constants in $10^{5}$ dynes $/ \mathrm{cm}$. are $f_{\nu}=9.6 f_{d}=2.87$

$$
f_{d \alpha}=0.575 \quad f_{a}-f_{a a}=0.317 \quad f_{\beta}-f_{\beta \beta}=0.23 \quad f_{a a}+f_{\beta \beta}=0.0034
$$

In the case of $\mathrm{PCl}_{3}$ molecule Howard and Wilson (1934) have obtained a value of $2.1 \times 10^{5}$ dynes $/ \mathrm{cm}$. for the force constant $f_{d}$ and Goutousky and Liehr have obtained the value of $11.4 \times 10^{5}$ dynes $/ \mathrm{cm}$. for $f_{v}$ in OPF $_{3}$. When an attempt was made to calculate the force constant assuming the value of $f_{\mathrm{p}}$ from Goutousky and

Liehr reasonable values for the other force constants could not be obtained possibly because of the large difference between the values for the two molecules. The value of $f_{d}$ is not much different from that for $\mathrm{PCl}_{3}$ whereas between $\mathrm{PF}_{3}$ and $\mathrm{OPF}_{3}$ the values obtained by Goutousky and Liehr were $4.64 \times 10^{5}$ dynes $/ \mathrm{cm} .^{3}$ and $6 \cdot 18 \times 10^{5}$ dynes $/ \mathrm{cm}$. respectively. For the particular potential function assumed the two angle angle interaction terms could be obtained only in combination the value was found to be rather small.

In conclusion the author wishes to express his grateful thanks to Prof. R. S. Krishnan for his kind encouragement and to Dr. P. S. Narayanan for valuable discussions.


