J. Indian Inst. Sci., July-Aug. 1993, 73, 347-363. ² Indian Institute of Science.

Nature of the spectrum associated with a matrix differential operator

S. DASKANUNGO¹ AND B. RAY PALADHI²

¹Department of Mathematics, Ramthakur College, Agartala 799 003, India.

²Department of Mathematics, J.C.C. College, Calcutta 700 033, India.

Abstract

The asymptotic solution of a self-adjoint matrix differential system are obtained by means of an integral transformation, using Titchmarsh's complex variable method. This leads to finding the nature of the spectrum of the differential system.

Key words and phrases: Matrix differential operator, eigenvalue parameter, self-adjoint, integral equation, Conte and Sangren's lemma, spectrum, L²-solution.

1. Introduction

We consider the differential system

$$(L - \lambda F) U = 0 \tag{1}$$

where,

$$L = \begin{pmatrix} \frac{d}{dx} \left(p_0(x) \frac{d}{dx} \right) + p_1(x) & r(x) \\ r(x) & \frac{d}{dx} \left(q_0(x) \frac{d}{dx} \right) + q_1(x) \end{pmatrix}$$

is a matrix differential operator, F, a 2×2 matrix and $U \cong \{u, v\}$, λ , a scalar. The system (eqn 1) is equivalent to

$$p_0 \frac{d^2 u}{dx^2} + p'_0 \frac{du}{dx} + p_1 u + rv - \lambda (F_{11} u + F_{12} v) = 0,$$

S. DASKANUNGO AND B. RAY PALADHI

$$q_0 \frac{d^2 \nu}{dx^2} + q_0' \frac{d\nu}{dx} + q_1 \nu + ru - \lambda (F_{21}u + F_{22}\nu) \approx 0$$
⁽²⁾

 $0 \le x < \infty$, where λ is an eigenvalue parameter; $p_0, p'_0, q_0, q'_0, \epsilon c[0, \infty)$; $p_1, q_1, r \to \infty$, as $x \to \infty$; $p_1, q_1, rcc[0,b]$ for all $b \rangle 0$; $p_0, q_0 \rangle 0$ in $[0,\infty]$; F is PDC $[0, \infty]$, along with the boundary conditions at x = 0, viz.,

i)
$$p_0(0) [a_{j1}u(0) + a_{j2}u'(0)] + q_0(0)[a_{j3}v(0) + a_{j4}v'(0)] = 0, (j = 1, 2).$$
 (3)

- ii) $a_{ik}(j = 1,2; k = 1,2,3,4)$ are real-valued constants.
- iii) the set $\{a_{1k}; k = 1,2,3,4\}$ is linearly independent of the set $\{a_{2k}; k = 1,2,3,4\}$.

iv)
$$q_0(0) (a_{14}a_{23} - a_{24}a_{13}) + p_0(0)(a_{12}a_{21} - a_{14}a_{22}) = 0$$
 (4)

The relation (4) ensures the self-adjointness of the system (2-3).

In the present paper, we study the nature of the spectrum of the self-adjoint differential system²⁻³, the system being considered earlier by Bhagat¹ in solving some other eigenvalue problem, and obtain some generalisation of the results in Paladhi².

The results and notations of Bhagat^{1,3}, Chakraborty⁵, Paladhi² and Titchmarsh⁵ are followed.

2. Notations

We use the following notations:

$$\begin{aligned} z(x) &= (\lambda - p_1(x)) \ (\lambda - q_1(x)) \\ \begin{pmatrix} u(x) \\ v(x) \end{pmatrix} &\equiv \{u(x), v(x)\} \equiv \{u, v\} \\ K(p_0 p'_0) &= [(p_0 z(x)^{-1/4})' + p_0 z(x)^{1/2} (z(x)^{1/4})'] [z(x)^{-3/4} p'_0 p_0^{-1} + (z(x)^{1/4})' z(x)^{-1}] \\ & (2 - 2 - 2) \end{aligned}$$

$$(Fij) = \begin{pmatrix} F_{11} & F_{12} \\ F_{21} & F_{22} \end{pmatrix}$$

F = (Fij) is PDC $[0, \infty]$ means that F is positive definite and continuous over the interval $[0, \infty]$.

$$\begin{aligned} (\alpha, \ \beta) &= \alpha_1 \overline{\beta}_1 + \alpha_2 \overline{\beta}_2 \text{ for any two vectors } \alpha = \{\alpha_1, \ \alpha_2\}, \ \beta = \{\beta_1, \ \beta_2\} \\ T(x) &= -i p_0^{-0} p_0' z(x)^{-3/4} \\ T_1(x) &= \{1/4 \ p_0 z(x)^{-5/4} \ z'(x) + p_0'(x)(z(x)^{-1/4})'(z(x)^{-1/4} \ p_0^{-1}) \end{aligned}$$

$$\begin{split} \xi_{1}(t) &= -z(t)^{1/2} + p_{0}^{-1}(p_{1} - \lambda F_{11}) \ z(t)^{-1/2} - T'(t) - iT_{I}(t) \\ &+ i \ p_{0}^{-2}(p_{0}z(t)^{1/4})' \ z(t)^{-1}p_{0}' \\ \eta_{1}(t) &= \{(r - \lambda F_{12}) \ q_{0}^{-1}z(t)^{-1/2} \} \\ l(t) &= \{\xi_{1}(t), \ \eta_{1}(t)\}, \ \Omega(t) &= \{\eta(t), \ \xi(t)\} \\ \zeta_{1}(t) &= \{-z(t)^{1/2} + q_{0}^{-1}(q_{1} - \lambda F_{22}) \ z(t)^{-1/2} - S'(t) - iS_{1}(t) \\ &+ i q_{0}^{-2} \ (q_{0}z(t)^{1/4}) \ z(t)^{-1}q_{0}' \}. \end{split}$$

This may be noted that S(t), $S_1(t)$ denote the same expressions for T(t), $T_1(t)$, respectively, p_0 and p'_0 being replaced by q_0 , q'_0 , respectively.

3. Downgrading of the coefficients of the differential system (2)

By using the transformation

$$\begin{split} \xi(x) &= i \int_{0}^{x} z(t)^{1/2} dt. \\ \begin{pmatrix} \eta(x) \\ \xi(x) \end{pmatrix} &\equiv \{\eta, \xi\} = i z(x)^{1/4} \{ p_0 u(x), q_0 \nu(x) \} \end{split}$$
(5)

(compare Paladhi²), we obtain

$$\begin{aligned} \frac{d\eta}{d\xi} &= [z(x)^{-1/4} \left(p_0 u' + p'_0 u \right) + \{z(x)^{1/4}\}' z(x)^{-1/2} p_0 u(x)] \\ \frac{d^2 \eta}{d^2 \xi^2} &= \frac{d}{d\xi} \left(\frac{d\eta}{d\xi} \right) = \frac{d}{dx} \left(\frac{d\eta}{d\xi} \right) \frac{dx}{d\xi} \\ &= -\frac{1}{p_0} \left[z(x)^{-1} (\lambda F_{11} + p''_0 - p_1) + \{ (z(x)^{-1/4})' (z(x)^{1/4} \times z(x)^{-1/2} - p'_0 z(x)^{-1/4} + p_0 z(x)^{-3/4} \{ z(x)^{-1/2} (z(x)^{1/4}) - (p_0 z(x)^{-1/4} + p_0 z(x)^{-1/2} (z(x)^{1/4})' \} \left\{ z(x)^{-3/4} \frac{p_0'}{p_0} + (z(x)^{1/4})' z(x)^{-1} \right\} \eta \\ &- \frac{1}{q_0} \left[z(x)^{-1} (\lambda F_{12} - r) \right] \zeta - iz(x)^{-1/4} \left\{ (z(x)^{-1/4} + (z(x)^{1/4})' \right\} \end{aligned}$$

$$z(x)^{-1/2} + \frac{p'_0}{p_0} z(x)^{-1/4} \frac{d\eta}{d\xi}$$

with similar expression for $\frac{d^2\xi}{d\xi^2}$.

Thus, (2) transforms to

In (6), the coefficient of η

$$\frac{d^{2}\eta}{d\xi^{2}} = -\frac{1}{p_{0}} \left[z(x)^{-1} (\lambda F_{11} + p_{0}'' - p_{1}) + \{ (z(x)^{-1/4})' + (z(x)^{1/4})' z(x)^{-1/2} \} \right]$$

$$p_{0}' z(x)^{-3/4} + p_{0} z(x)^{-3/4} + p_{0} z(x)^{3/4} \{ z(x)^{-1/2} (z(x)^{1/4})' \}'$$

$$- (p_{0}p_{0}') \eta - \frac{1}{q_{0}} \left[z(x)^{-1} (\lambda F_{12} - r) \right] \} - iz(x)^{-1/4} \{ (z(x)^{-1/4})' + (z(x)^{1/4})' z(x)^{-1/2} + \frac{p_{0}'}{p_{0}} z(x)^{-1/4} \} \frac{d\eta}{d\xi}$$
(6)

and

$$\frac{d^{2}\zeta}{d\xi^{2}} = -\frac{1}{q_{0}} \left[z(x)^{-1} (\lambda F_{22} + q_{0}'' - q_{1}) + \left\{ (z(x)^{-1/4})' + (z(x)^{1/4})' z(x)^{-1/2} \right\} \right]$$

$$q_{0}' z(x)^{-3/4} + q_{0} z(x)^{-3/4} \left\{ z(x)^{-1/2} (z(x)^{1/4})' \right\}' - K(q_{0}q_{0}') \left\{ \zeta - \frac{1}{p_{0}} \left[z(x)^{-1} (\lambda F_{21} - r) \right] \eta - i z(x)^{-1/4} \left\{ (z(x)^{-1/4})' + (z(x)^{-1/4} + \frac{q_{0}'}{q_{0}} \right\} \right]$$

$$z(x)^{-1/4} \frac{d\zeta}{d\xi}.$$
(7)

$$= -\frac{1}{p_0} \left[z(x)^{-1} (\lambda F_{11} + p_0'' - p_1 + \{ (z(x)^{-1/4})' + (z)x)^{1/4})' \times z(x)^{-1/2} \} \right. \\ \left. p_0' z(x)^{-3/4} + p_0 z(x)^{-3/4} \left\{ z(x)^{-1/2} (z(x)^{1/4})' \right\}' - K(p_0 p_0') \right]$$

→ a finite number under suitable conditions on p_0 , p''_0 , p_1 , q_1 , F_{11} and the coefficient of ξ and $\frac{d\eta}{d\xi}$ are 0(1) as $x \to \infty$. Similarly, in (7), the coefficient of η , ξ , $\frac{d\zeta}{d\xi}$ is 0(1) as $x \to \infty$.

4. Derivation of an integral solution of the system (6,7)

Let $p(\mathbf{x}) = z(\mathbf{x})^{1/4} H(\mathbf{x}), \ H(\mathbf{x}) = \{H_1(\mathbf{x}), \ H_2(\mathbf{x})\}$

where
$$H_1(x) = \frac{d}{dx} \left[z(x)^{-1/2} \frac{d\eta}{dx} \right] - i \ z(x)^{-1/4} \frac{d}{dx} \left(p_0 \frac{du}{dx} \right)$$
(8)

$$H_{2}(x) = \frac{d}{dx} \left[z(x)^{-1/2} \frac{d\zeta}{dx} \right] - i \ z \ (x)^{-1/4} \frac{d}{dx} \left(q_{0} \frac{d\nu}{dx} \right)$$
(9)

Then,

$$p(x) = z(x)^{1/4} \begin{pmatrix} H_1(x) \\ H_2(x) \end{pmatrix}$$
where $z(x)^{1/4}H_1(x) = z(x)^{1/4} \frac{d}{dx} \left[z(x)^{-1/2} \frac{d\eta}{dx} \right] - i \frac{d}{dx} \left(p_0 \frac{du}{dx} \right)$

$$= i z(x)^{1/4} \left[\left\{ \frac{p_0}{4} z(x)^{-5/4} z'(x) + p'_0 z(x)^{-1/4} \right\}' u(x) + (p'_0 z(x)^{-1/4})u'(x).$$
(10)

Similarly,

$$z(x)^{1/4} H_2(x) = iz(x)^{1/4} \left[\left\{ \frac{q_0}{4} z(x)^{-5/4} z'(x) + q'_0 z(x)^{-1/4} \right\}' \nu(x) + (q'_0 z(x)^{-1/4}) \nu'(x). \right]$$
(11)

We have,
$$\int_{0}^{x} \sin (\xi(x) - \xi(t)) H_{1}(t) dt$$

$$= \int_{0}^{x} \left[\frac{d}{dt} \left\{ z(t)^{-1/2} \frac{d\eta}{dt} \right\} + i \left\{ p_{1}u + rv - \lambda(F_{11}u + F_{12}v) \right\} z(t)^{-1/4} \right]$$

$$\times \sin (\xi(x) - \xi(t)) dt \text{ by (8 and 2)}$$

$$= -\sin \xi(x) \lambda^{-1} \eta'(0) + i\eta(x) - i\eta(0) \cos \xi (x) - \int_{0}^{x} \sin(\xi(x) - \xi(t)) z(t)^{1/2}$$

$$\times \eta(t) dt + \int_{0}^{x} \sin (\xi(x) - \xi(t)) \left\{ \frac{p_{1}}{p_{0}} \eta(t) + \frac{r}{q_{0}} \xi(t) - \lambda \left(F_{11} \frac{\eta(t)}{p_{0}} + F_{12} \frac{\zeta(t)}{q_{0}} \right) \right\} z(t)^{-1/2} dt.$$
(12)
Also by (10) we have

Also by (10) we have

$$\int_{0}^{x} \sin (\xi(x) - \xi(t)) H_{1}(t) dt.$$

$$= i \int_{0}^{x} \operatorname{Sin} \left(\xi(x) - \xi(t)\right) \left[\left\{ \frac{p_0}{4} z(t)^{-5/4} z'(t) + p'_0 z(t)^{-1/4} \right\}' u(t) + \left\{ p'_0 z(t)^{-1/2} \right\} u'(t) dt. \right]$$
(13)

From (12-13), we obtain

$$\begin{aligned} \eta(x) &= \eta(0) \, \cos \, \xi(x) - i \eta^{-1} \eta' \, (0) \, \sin \, \xi(x) + i \int_{0}^{x} \sin \, \left(\xi(x) - \xi(t) \right) \\ &\times \left[\left\{ - z(t)^{1/2} + p_{0}^{-1} \, \left(p_{1} - \lambda F_{11} \right) z(t)^{-1/2} \right\} \eta(t) + \left(r - \lambda F_{12} \right) \, q_{0}^{-1} \\ &z(t)^{-1/2} \xi(t) + K(t, \, \lambda) \right] \, dt, \end{aligned}$$
(14)

where,
$$K(t, \lambda) = \left\{ \frac{p_0(t)}{4} z(t)^{-5/4} z'(t) + p'_0(t)z(t)^{-1/4} \right\}' u(t) + p'_0 z(t)^{-1/2} u'(t).$$
(15)

We have from (5)

$$\begin{split} \eta(x) &= i \; z(x)^{1/4} \; p_0 \; (x) \; u(x) \\ \eta'(x) &= i [(z(x)^{1/4} \; p_0')' \; u(x) \; + \; z(x)^{1/4} \; p_0(x) u'(x)] \end{split}$$

so that

$$u'(x) = i \left[-z(x)^{-1/4} p_0^{-1} \eta'(x) + (z(x)^{1/4} p_0)'(z(x)^{-1/2} p_0^{-2} \eta(x)) \right]$$
(16)

It then follows from (15-16) that

$$K(t, \lambda) = T(t) \eta'(t) - i \left\{ \frac{p_0}{4} z(t)^{-5/4} z'(t) + p'_0 z(t)^{-1/4} \right\}'$$

$$(z(t)^{-1/4} p_0^{-1}) \eta(t) + i (p_0 z(t)^{1/4})' p'_0 p_0^{-2} z(t)^{-1}) \eta(t)$$

$$= T(t) \eta'(t) - iT_1(t) \eta(t) + i p_0^{-2} (p_0 z(t)^{1/4})' z(t)^{-1} p'_0 \eta(t)$$
(17)

Now, $\int_{0}^{x} \sin (\xi (x) - \xi(t)) T(t) \eta' (t) dt$ $= - T(0) \eta(0) \sin \xi(x) - \int_{0}^{x} [\sin (\xi(x) - \xi(t)) T'(t) - i \cos (\xi (x) - \xi(t)) z(t)^{1/2} T(t)] \eta(t) dt.$ (18)

It follows then from (14, 15, 17 and 18) that

$$\begin{aligned} \eta(x) &= \eta(0) \cos \xi(x) + (-i\lambda^{-1}\eta'(0) - i \ T(0)\eta(0) \sin \xi(x) \\ &+ i \int_{0}^{x} \left[\sin \left(\xi \ (x) - \xi(t) \right) \right] \left\{ -z(t)^{1/2} + \right. \\ &+ p_{0}^{-1} \ (p_{1} - \lambda F_{11}) z(t)^{-1/2} - T'(t) - iT_{1}(t) + ip_{0}^{-2} \ (p_{0}z(t)^{-1/4})' \\ &- z(t)^{-1} \ p_{0}' \right\} \eta(t) + (r - \lambda F_{12}) \ q_{0}^{-1} \ z(t)^{-1/2} \right\} \ \zeta(t) \ dt \\ &- \int_{0}^{x} \cos \left(\xi(x) - \xi(t) \right) \ z(t)^{1/2} \ T(t) \ \eta(t) \ dt. \end{aligned}$$
(19)

Proceeding in the same way, taking (9) and (11) into consideration, we obtain

$$\begin{aligned} \zeta'(x) &= \zeta(0) \cos \zeta(x) + (-i \lambda^{-1} \zeta'(0) - iS(0)\zeta(0)) \sin \xi(x) \\ &+ i \int_{0}^{x} \sin \left(\xi(x) - \xi(t)\right) \left[\left\{ -z \left(t\right)^{1/2} + q_{0}^{-1} \left(q_{1} - \lambda F_{22}\right) z(t)^{-1/2} \right. \\ &- S'(t) - iS_{1}(t) + iq_{0}^{-2}(q_{0}z(t)^{1/4})' z(t)^{-1} q_{0}' \right\} \zeta(t) + \left\{ (r - \lambda F_{12}) \right. \\ &\times p_{0}^{-1} z(t) z(t)^{-1/2} \eta(t) \right] dt \\ &- \int_{0}^{x} \cos \left(\xi(x) - \zeta(t)\right) z(t)^{1/2} S(t) \xi(t) dt \end{aligned}$$
(20)

where S(t) and $S_1(t)$ can be obtained from T(t) and T(t), respectively, replacing p_0 by q_0 in the expressions for T(t), $T_1(t)$ (see § 2).

Equations (20) can be written in the form

$$\Omega(x) = \begin{pmatrix} \eta & (0) \\ \xi & (0) \end{pmatrix} \cos \xi (x) + i \begin{pmatrix} -\lambda^{-1} \eta' & (0) - T(0) & \eta(0) \\ -\lambda^{-1} \zeta' & (0) - S(0) & \zeta(0) \end{pmatrix} \sin \xi (x) + i \int_{0}^{x} \sin \left(\xi(x) - \xi(t)\right) \begin{pmatrix} \xi_{1} & \eta_{1} \\ \eta_{1} & \zeta_{1} \end{pmatrix} \begin{pmatrix} \eta & (t) \\ \zeta & (t) \end{pmatrix} dt$$
(21)
$$- i \int_{0}^{x} \cos \left(\xi(x) - \xi(t)\right) z(t)^{1/2} \begin{pmatrix} T(t) & \eta(t) \\ S(t) & \zeta(t) \end{pmatrix} dt.$$
(21)

5. Some propositions

Proposition 1: If the coefficients p_0 , q_0 , p_1 , q_1 , r(x) in the differential system (2) satisfy the following conditions, viz.,

- (i) $p_1(x), q_1(x) \to \infty$ as $x \to \infty$, *i.e.*, $p_1(x), q_1(x) > Q(x)$ whenever $Q(x) \ge \delta > 0$; $x \ge 0$; $F_{12} q_0^{-1}, r(x)q_0^{-1}, \in L[O, \infty)$; $p_0^{-1} \{p_1, F_{11}\}, \{q_1 \ F_{22}\} q_0^{-1}$ have all its elements square integrable over $[0, \infty)$.
- (ii) $p'_1(x), q'_1(x) \ge 0$
- (iii) $p'_1(x) = 0[p_1(x)]^c$, $q_1(x) = 0[q_1(x)]^c$, $0 \times c < 7/4$
- (iv) $p_1''(x)$, $q_1''(x)$ maintain their signs

(v)
$$\frac{1}{Q(x)} \epsilon L[0, \infty)$$

(vi) $p_0' p_0^{-1}$, $q_0 q_0^{-1}$, $(p_0^{-1} p_0')'$, $(q_0^{-1} q_0)' \in L$ [0, ∞) and $p_0 p_0'$, $q_0 q_1'$, $p_0'' p_0^{-1}$, $q_0'' q_0^{-1} = 0$ (1) (compare Titchmarsh⁵, (Ch. V, p. 121) and Paladhi² p. 448), then the integrals

$$\int_{0}^{\infty} \left[\left| \xi(t) \right| + \left| \eta_{1}(t) \right| \right] \left| z(t)^{-1} \right| dt, \int_{0}^{\infty} \left[\left| \eta_{1}(t) \right| + \left| \zeta_{1}(t) \right| \right] \left| z(t)^{-1} \right| dt$$

and

$$\int_{0}^{\infty} |T(t)| |z(t)^{-1/2} | dt, \quad \int_{0}^{\infty} |S(t)| |z(t)^{-1/2} | dt \text{ are uniformly convergent}$$

with respect to λ over $|\lambda - p_1(x)|$, $|\lambda - q_1(x) \ge \delta_1 > 0$ for $0 \le x < \infty$. Proof. We have,

$$\int_{A}^{x} \left[\left| \xi_{1}(t) \right| + \left| \eta_{1}(t) \right| \right] |z(t)^{-1}| dt$$

$$\leq \int_{A}^{x} |z(t)^{-1/2}| dt + \int_{A}^{x} |p_{0}^{-1}(p_{1} - \lambda F_{11})| |z(t)^{-3/2}| dt + \int_{A}^{x} |T'(t)| |z(t)^{-1}| dt$$

$$+ \int_{A}^{x} |T_{I}(t)| |z(t)^{-1}| dt + \int_{A}^{x} |p_{0}^{-2}(p_{0}z(t)^{1/4}p_{0}')| |z(t)^{-1}| dt$$

$$+ \int_{A}^{x} |(r - \lambda F_{12})|q_{0}^{-1}| |z(t)^{-3/2}| dt = I_{11} + I_{12} + I_{13} + I_{14} + I_{15} + I_{16}, \text{ say.}$$

Now,
$$I_{11} = \int_{A}^{x} |z(t)^{-1/2}| dt$$

$$\leq \int_{A}^{x} Q(t)^{-1} dt, \text{ since } |z(t)| > Q^{2}(t) \text{ for } t > A$$

$$\leq \int_{A}^{x} |p_{0}^{-1}(p_{1} - \lambda F_{11})| Q(t)^{-3} dt$$

$$\leq \left[\int_{A}^{x} |p_{0}^{-1}(p_{1} - \lambda F_{11})|^{2} dt\right]^{12} \left[\int_{A}^{x} Q(t)^{-6} dt\right]^{12}, \text{ by Schwarz}$$
inequality.

$$= 0(1) \text{ as } x \to \infty, \text{ by conditions (i) and (v).}$$

$$I_{13} = \int_{A}^{x} |T'(t)| ||z(t)^{-1} dt$$

$$\leq \int_{A}^{x} |3/4 p_{0}^{-1} p_{0}| \{(\lambda - p_{1}(t)) q_{1}(t) + (\lambda - q_{1}(t))p_{1}(t)\}|$$

$$|z(t)^{-11/4}| dt + \int_{A}^{x} |(p_{0}^{-1} p_{0})'| ||z(t)^{-7/4}| dt$$

$$= 0 (1), \text{ provided } 0 < c < 11/4.$$

$$I_{14} = \int_{A}^{x} |T_{1}(t)| ||z(t)^{-1}| dt$$

$$= 0 \int_{A}^{x} [|z(t)|^{-5/2}||z'(t)| + |p_{0}(t)p_{0}^{-1}|| ||z(t)^{-5/2}|||z'(t)| ||t|$$

$$+ |z(t)^{-7/2}|||z'(t)^{2}|| + |p_{0}'(t)||z(t)^{-3/2} p_{0}^{-1}||]dt$$

$$= 0 \int_{A}^{x} [p_{1}^{-5/2} q_{1}^{-5/2} (p_{1}^{c} q_{1}^{c} + p_{1} q_{1}'' + q_{1} p_{1}'') + p_{1}^{-5/2} q_{1}^{-5/2} (p_{1} q_{1}^{c} + q_{1} p_{1}^{c})]$$

ACCREME.

$$+ p_1^{-7/2} q_1^{-7/2} (p_1^2 q_1^{2c} + p_1^{2c} q_1^2) + |p_0^{c}(t)p_0^{-1}| (p_1^{-3/2} q_1^{-3/2}) dt$$

$$= 0 (1), \text{ as } x \to \infty \text{ if } 0 < c < 7/4.$$

$$I_{15} = \int_A^x |p_0^{-2} (p_0 z(t)^{1/4})' p_0'| |z(t)^{-1}| dt$$

$$= 0 \int_A^x p_0'^2 p_0^{-2} p_1^{-3/4} q_1^{-3/4} + \frac{p_0 p_0'}{4} p_1^{-7/4} q_1^{-7/4} (p_1 q_1^c + q_1 p_1^c) dt$$

$$= 0(1), \ 0 < c < 7/4, \text{ by conditions (v) and (vi).}$$

$$I_{16} = \int_A^x |(r - F_{12}) q_0^{-1}| |z(t)^{-3/2}| dt \text{ as } x \to \nu$$

$$= 0(1), \ \text{as } x \to \infty, \text{ since } rq_0^{-1}, \ F_{12} q_0^{-1} \in [L[0,\infty) \text{ by condition (i).}$$

Thus, the integral

$$\int_{0}^{\infty} |z(t)^{-1}| (|\xi_{1}(t)| + |\eta_{1}(t)|) dt < \infty$$

over $|\lambda - p_{1}(x)|, |\lambda - q_{1}(x)| \ge \delta_{1} > 0$ for $0 \le x < \infty$

Similar results hold for the other integral

$$\int_{0}^{\infty} |z(t)^{-1}| (|\eta_{1}(t)| + |\zeta_{1}(t)|) dt.$$
Now,

$$\int_{0}^{\infty} |T(t)| |z(t)^{-1/2}| dt$$

$$= 0 \int_{0}^{\infty} |p_{0}^{-1}p_{0}'| p_{1}^{-5/4} q_{1}^{-5/4} dt = 0 \Big[\int_{0}^{\infty} Q(t)^{-5/2} dt \Big] = 0(1)$$
as $x \to \infty$, by condition (vi).
Similarly,

$$\int_{0}^{\infty} |s(t)| |z(t)^{-1/2}| dt = 0(1) \text{ as } x \to \infty, \text{ by condition (vi)}.$$

Thus, the proposition is established.

Proposition 2: If $Im \lambda > 0$, $0 < \arg \lambda < \pi$, then $\exp[i \xi(x)] \to \infty$ as $x \to \infty$ [see Paladhi², Lemma II].

6. Estimates for $\eta(x)$, $\zeta(x)$

Setting
$$\Omega^{1}(\mathbf{x}) = \{\eta^{\dagger}(\mathbf{x}), \zeta^{\dagger}(\mathbf{x})\} = \exp(-i\xi(\mathbf{x}))z(\mathbf{x})\Omega(\mathbf{x})$$
 (22)
in (19) and (20), we have
 $\eta(\mathbf{x}) = 1/2\eta(0)e^{-i\xi(\mathbf{x})}[1+e^{-2i\xi(\mathbf{x})}]+i(-\lambda^{-1}\eta'(0)-T(0)\eta(0)) \times \\ \times \frac{1}{2i}e^{i\xi(\mathbf{x})}[1-e^{-2i\xi(\mathbf{x})}]+i\int_{0}^{x} (\xi_{1}(t)\eta(t)+\eta_{1}(t)\zeta(t)) \times \\ \times e^{i(\xi(\mathbf{x})-\xi(t))}\frac{1}{2i}\left[1+e^{-2i(\xi(\mathbf{x})-\xi(t))}\right]dt \\ -\frac{1}{2}\int_{0}^{x}e^{i(\xi(\mathbf{x})-\xi(t))}\left[1+e^{-2i(\xi(\mathbf{x})-\xi(t))}\right]z(t)^{1/2}T(t)\eta(t) dt \\ \text{or, } \eta^{+}(\mathbf{x}) = \frac{1}{2}\eta(0)\left[1+e^{-2i\xi(\mathbf{x})}\right]z(\mathbf{x}) + i(-\lambda^{-1}\eta(0)-T(0)\eta(0))z(\mathbf{x}) \\ \times \frac{1}{2i}\left[1-e^{-2i\xi(\mathbf{x})}\right] + iz(\mathbf{x})\int_{0}^{x}(\xi_{1}(t)\eta^{+}(t)+\eta_{1}(t)\xi^{+}(t)) \\ \times z(t)^{-1}\frac{1}{2i}\left[1-e^{-2i(\xi(\mathbf{x})-\xi(t))}\right]dt \\ -\frac{z(\mathbf{x})}{2}\int_{0}^{x}\eta^{+}(t)z(t)^{-1/2}[1+e^{-2i(\xi(\mathbf{x})-\xi(t))}]T(t) dt. \\ \text{Therefore, } \left|\frac{\eta^{+}(\mathbf{x})}{z(\mathbf{x})}\right| \leq \left|\eta(0)\right| + \left|-\lambda^{-1}\eta'(0)-T(0)\eta(0)\right| + \\ \int_{0}^{x}\left(|\xi_{1}(t)| |\eta^{+}(t)| + |\eta_{1}(t)\zeta^{+}(t)|\right)|z(t)^{-1}|dt + \int_{0}^{x}|T(t)z(t)^{-1/2}| \\ |\eta^{+}(t)| dt$ (23)

Similarly,

$$\begin{aligned} \frac{|\zeta^{+}(x)|}{z(x)} &|\leq |\zeta(0)| + |-\lambda^{-1}\zeta'(0) - S(0)\zeta(0)| + \int_{0}^{x} (|\eta_{1}(t)|| |\eta^{+}(t)| + \\ + |\zeta_{1}(t)|| |\zeta^{+}(t)||) |z(t)^{-1}| dt + \int_{0}^{x} |S(t)z(t)^{-1/2}|| |\zeta^{+}(t)| dt. \end{aligned}$$
(24)

Now, (23) and (24) can be written in the form

$$\left|\frac{\eta^+(x)}{z(x)}\right| \leq A + \int_0^x \left[\left| \frac{\xi_1^+(t)}{z(t)} \right| \right| \eta^+(t) \left| + \left| \frac{\eta_1(t)}{z(t)} \right| \left| \zeta^+(t) \right| \right] dt$$

and
$$\left|\frac{\zeta^{+}(x)}{z(x)}\right| \leq A + \int_{0}^{x} \left[\left|\frac{\eta_{1}(t)}{z(t)}\right| + \eta^{+}(t)\right| + \left|\frac{\zeta^{+}(t)}{z(t)}\right| + \left|\zeta^{+}(t)\right| \right]_{dt}$$
(25)

where
$$A = \max \left[1 - \lambda^{-1} \eta'(0) - T(0) \eta(0) \right] + \eta(0) + \eta(0) + \eta(0) - S(0)\zeta(0) + \zeta(0) + \zeta(0$$

and
$$|\xi_1^{\dagger}(t)| = |\xi_1(t)| + |T(t) z(t)^{1/2}|$$

 $|\zeta_1^{\dagger}(t)| = |\zeta_1(t)| + |S(t) z(t)^{1/2}|.$ (26)

We can take
$$|\xi_1^+(t)| = \max[|\xi_1^+(t)|, |\zeta_1^+(t)|] = M_1(t)$$
 say, (27)

for otherwise the same lines of argument will follow with

$$\left| \zeta_{1}^{+}(t) \right| = \max\left[\left| \xi_{1}^{+}(t) \right|, \left| \zeta_{1}^{+}(t) \right| \right]. \text{ It follows then from (25) that} \\ \left| \frac{\eta^{+}(x)}{z(x)} \right|, \left| \frac{\zeta^{+}(x)}{z(x)} \right| \leq A + \int_{0}^{x} (\alpha, \beta) dt$$
(28)

where $(\alpha, \beta) = \max \left[\left(\{ M_{I}(t), | \eta_{1}(t) | \}, \beta \right), \left(\{ | \eta_{1}(t) |, M_{1}(t) \}, \beta \right) \right]$ and $\beta = \left\{ \left| \frac{\eta^{+}(t)}{z(t)} \right|, \left| \frac{\zeta^{+}(t)}{z(t)} \right| \right\}$; if $\alpha = \{\alpha_{1}, \alpha_{2}\}$, then $\alpha_{1} + \alpha_{2} = M_{1}(t) + |\eta_{1}(t)|$ (29)

Thus,
$$\left| \frac{\eta^{+}(x)}{z(x)} \right|$$
, $\left| \frac{\zeta^{+}(x)}{z(x)} \right| \leq A + \int_{0}^{x} \left[\alpha_{1} \left| \frac{\eta^{+}(t)}{z(t)} \right| + \alpha_{2} \left| \frac{\zeta^{+}(t)}{z(t)} \right| \right] dt$

$$(30)$$

$$\leq A \exp \left[\int_{0}^{x} (\alpha_{1}+\alpha_{2}) dt, \text{ by Conte and Sangren}^{7}\right]$$
$$\leq A \exp \left[|z(x)| \int_{0}^{\infty} (\alpha_{1}^{+}(t) + \alpha_{2}^{+}(t)) dt\right], \text{ where } \alpha_{j}^{+}(t) = \alpha_{j}(t) |z(t)|^{-1}.$$

It follows by definition of $\Omega^+(x)$ that

$$|\eta(x) \exp(-i\xi(x))|, |\zeta(x) \exp(-i\xi(x))|$$

or
$$\leq A \exp \int_{0} \left\lfloor \frac{\lambda \mathcal{U}_{1}(t)}{|z(t)|} + \frac{|\eta_{1}(t)|}{|z(t)|} \right\rfloor |z(x)| dt$$

 $|\eta(x)|, |\zeta(x)| \leq A |\exp(i\xi(x))|$

$$\times \exp\left[|z(x)| \int_{0}^{\infty} \frac{M_{1}(t) + |\eta_{1}(t)|}{|z(t)|} dt \right].$$
(31)

Now,
$$\int_{0}^{\infty} \frac{M_{1}(t) + |\eta_{1}(t)|}{|z(t)|} dt$$

is either equal to
$$\int_{0}^{\infty} \frac{|\xi_{1}^{+}(t)| + |\eta_{1}(t)| dt}{|z(t)|}$$

$$= \int_{0}^{\infty} \frac{|\xi_{1}(t)| + |\eta_{1}(t)|}{|z(t)|} dt + \int_{0}^{\infty} |T(t)|z(t)^{-1/2}| dt$$

or equal to
$$\int_{0}^{\infty} \frac{|\zeta_{1}^{+}(t)| + |\eta_{1}(t)|}{|z(t)|} dt + \int_{0}^{\infty} |s(t)|| |z(t)^{-1/2}| dt.$$

The two integrals above in the R.H.S. are convergent by Proposition 1. Therefore, we obtain from (31)

$$|\eta(x)|, |\zeta(x)| \leq A \exp(i\xi(x))) \exp[k | z(x) |]$$

= 0[| exp(i\xi(x)) | | exp(K p_i(x)q_i(x)) |]. (32)

2.7 Derivation of asymptotic expressions for $\{\eta(x),\zeta(x)\} = \Omega(x)$ From (21), we have

$$\begin{aligned} \eta(x) &= \exp(i\xi(x) \left[\frac{1}{2} (1 + \exp(-2i\xi(x))) \eta(0) + \frac{1}{2} (1 - \exp(-2i\xi(x))) (-\lambda^{-1}\eta^{1}(0) - T(0) - \eta(0) + \frac{1}{2} z(x) \exp(Kp_{1}(x)q_{1}(x)) \right] \\ &\times \int_{0}^{x} z(x)^{-1} \exp(-Kp_{1}(x)q_{1}(x)) \exp(-i\xi(t))(1 - \exp(-2i(\xi(x) - \xi(t)))) \\ &\times (\xi_{1}(t)\eta(t) + \eta_{1}(t)\zeta(t)) dt \\ &- \frac{1}{2} \exp(Kp_{1} q_{1}) z(x^{-1}) \int_{0}^{x} z(x^{-1}) \{1 + \exp(-2i(\xi(x) - \xi(t)))\} \\ &\times z(t)^{1/2} T(t) \eta(t) \exp(-K p_{1}(x)q_{1}(x) \exp(-i\xi(t)) dt]. \end{aligned}$$
(33)

Now,

$$|z(x)^{-1}\exp(-Kp_{1}(x)q_{1}(x)) \exp(-i\xi(t)) \{1-\exp(-2i\xi(x)-\xi(t))\}$$

$$\times (\xi_{1}(t) \eta(t) + \eta_{1}(t) \zeta(t)) |$$

$$\leq z(t^{-1}) \exp(-Kp_{1}(t)q_{1}(t) \exp(-i\xi(t)) (|\xi_{1}(t)| | \eta(t)| +$$

$$+ |\eta_{1}(t)| | \zeta(t)|)$$

$$= |\exp(-Kp_{1}(t)q_{1}(t)| |\exp(-i\xi(t)|$$

$$\times \left(\frac{|\xi_{1}(t)|}{|z(t)|} | \eta(t)| + \frac{|\eta_{1}(t)|}{|z(t)|} | \zeta(t)| \right)$$

$$\leq K_{I} \left(\frac{|\xi_{1}(t)|}{|z(t)|} + |\eta_{1}(t)| \right), \text{ by (32) when } K_{1} \text{ is a constant.}$$
(34)
Also, $|z(x)^{-1}\{1+\exp(-2i(\xi(x)-\xi(t))\} | z(t)^{1/2} T(t)\eta(t)\exp(-Kp_{1}(x)q_{1}(x))$

$$\times \exp(-i\xi(t))|$$

$$\leq |\{1+\exp(-2i(\xi(x)-\xi(t)))\}z(t)^{-1/2}| | T(t)\eta(t)| | \exp(-i\xi(t)|$$

$$\times |\exp(-Kp_{1}(t)q_{1}(t)|$$
But, $\int_{0}^{\infty} (|\xi_{1}(t)| + |\eta_{1}(t)|) | z(t)^{-1}| dt, \int_{0}^{\infty} |T(t)| z(t)^{-1/2}| dt \text{ are convergent by}$

But, $\int_{0} (|\xi_{1}(t)| + |\eta_{1}(t)|) |z(t)^{-1}| dt$, $\int_{0} |T(t) z(t)^{-1/2}| dt$ are convergent by Proposition 1.

Therefore, from (34) and (35), it follows that

$$\begin{split} L_1 &= \lim_{x \to \infty} \int_0^x + z(x)^{-1} \exp(-Kp_1(x)q_1(x))\exp(-i\xi(t)) \times \{1 + \exp(-2i(\xi(x) - \xi(t))\} \\ &\times (\xi_1(t) \ \eta(t) + \eta_1(t)\zeta(t)) \ dt, \text{ and} \\ L_2 &= \lim_{x \to \infty} \int_0^x -z(x)^{-1} \exp(-Kp_1(x)q_1(x)) \ \exp(-i\xi(t)) \\ &\times \{1 + \exp(-2i(\xi(x) - \xi(t)))\}z(t)^{1/2}T(t)\eta(t) \ dt \\ \text{are finite numbers.} \end{split}$$

Hence, it follows from (33) $\eta(x) \sim \frac{L^*}{2} z(x) \exp(Kp_1(x)q_1(x) + i\xi(x))$ where $L^* = L_1 + L_2$. (36)

Proceeding in a similar way, it follows that

$$\xi(x) \sim S^*/2 \ z(x) \ \exp(Kp_1(x)q_1(x) + i\xi(x)) \ \text{where} \ S^* = S_1 + S_2$$
(37)

with
$$S_1 = \lim_{x \to \infty} \int_0^x z(x)^{-1} \exp(-Kp_1(x)q_1(x) - i\xi(t)) \{1 - \exp(-2i(\xi(x) - \xi(t)))\}$$

 $\times (\eta_1 \eta(t) + \zeta_1 \zeta(t) dt$

and $S_2 = \lim_{x \to \infty} \int_0^x -z(x)^{-1} \exp(-Kp_1(x)q_1(x) - i\xi(t)) \{1 + \exp(-\xi(x) - \xi(t))\}$

$$\times z(t)^{1/2} S(t) \zeta(t) dt.$$

We take

$$U_{j}(\mathbf{x}) = \begin{pmatrix} U_{j1} \\ U_{j2} \end{pmatrix} = \{U_{j1}, U_{j2}\}, V_{j}(\mathbf{x}) = \begin{pmatrix} V_{j1} \\ V_{j2} \end{pmatrix}$$

where $\{U_{j1}, U_{j2}\} = U_j(x) = iz(x)^{1/4} \{p_0 x_j(x), q_0 y_j(x)\}$ $V_j(x) = iz(x)^{1/4} \{P_0 u_j(x), q_0 v_j(x)\}$ $= \{V_{j1}, V_{j2}\}, j = 1, 2.$

 $\phi_k = \{u_k, v_k\}, k = 1,2$ are two solution vectors of (2) satisfying (2-4) and $\theta_k = \{x_k, y_k\}$ are also two solution vectors of (2) such that

$$[\theta_1\theta_2] = 0$$
, $[\phi_j\theta_k] = \delta_{jk}(j,k = 1,2)$ and

$$[\phi_j\phi_k] = p_0 \qquad \begin{vmatrix} u_j & u_k \\ u_j' & u_k' \end{vmatrix} + q_0 \qquad \begin{vmatrix} v_j & v_k \\ v_j' & v_k' \end{vmatrix}$$

It follows from (36-37)

$$U_{j}(x,\lambda) \sim \frac{1}{2} z(x) \exp \left[K_{1}p_{1}(x)q_{1}(x)+i\xi(x)\right] L_{j}(\lambda)$$

$$V_{j}(x,\lambda) \sim \frac{1}{2} z(x) \exp \left[K_{1}p_{1}(x)q_{1}(x)+i\xi(x)\right] S_{j}(\lambda)$$
(38)

where $L_j(\lambda) \equiv \{L_{1j}(\lambda), L_{2j}(\lambda)\}, S_j(\lambda) \equiv \{S_{1j}(\lambda), S_{2j}(\lambda)\}.$

$$U_j(x,\lambda) = \{U_{j1}, V_{j2}\}, V_j(x,\lambda) = \{V_{j1}, V_{j2}\};$$

and L_{ij} are independent of x; $U_j(x) = iz(x)^{1/4} \{p_0 x_j(x), q_0 y_j(x)\}$

and
$$V_j(x) = i z(x)^{1/4} \{ p_0 u_j(x), q_0 v_j(x) \}$$

8. Main theorem

Theorem. The spectrum of the differential system (2-4) is a countably infinite point spectrum in $(-\infty,\infty)$ if all the conditions (i)-(vi) of Proposition 1 are satisfied by the potentials p_0 , q_0 , p_1 , q_1 , r(x).

Proof. We consider two solutions of (2-4), viz.,

$$\psi_{j} = \theta_{j}(x,\lambda) + \sum_{r=1}^{2} m_{r}(\lambda) \phi_{r}(x,\lambda) \qquad (j = 1,2)$$

$$= \begin{pmatrix} x_{j} + m_{j1} u_{1} + m_{j2} u_{2} \\ y_{j} + m_{j1} v_{1} + m_{j2} v_{2} \end{pmatrix}$$

$$= -iz(x)^{-1/4} \begin{pmatrix} p_{0}^{-1} U_{j1}(x) + m_{j1}p_{0}^{-1} V_{11} + m_{j2}p_{0}^{-1} V_{21} \\ q_{0}^{-1} U_{j2}(x) + m_{j1} q_{0}^{-1} V_{12} + m_{j2}q_{0}^{-1}V_{22} \end{pmatrix}$$

Therefore,

$$\begin{split} \psi_{j} &\sim (-\frac{1}{2}i) \ z(x)^{3/4} \exp[K_{1}p_{1}(x)q_{1}(x)+i\xi(x)] \left(\begin{matrix} p_{0}^{-1}[L_{1j} + m_{j1}S_{11} + m_{j2}S_{12}] \\ q_{0}^{-1}[L_{2j} + m_{j1} S_{21} + m_{j2}S_{22}] \end{matrix} \right) \\ &= + \frac{1}{2c} \ z(x)^{3/4} \ \exp[K_{1}p_{1}(x)q_{1}(x) + i\xi(x)] \ W_{j}(\lambda, \ p_{0}^{-1}, \ q_{0}^{-1}) \\ (j = 1, 2). \end{split}$$

$$(39)$$

But for non-real values of λ , we have in the singular case $0 \le x < \infty$ at least two linearly independent solutions of (2) say $W_r(x, \lambda)$ (r = 1,2) such that

$$W_r(x,\lambda) \in L^2[0, \infty)$$
 (Bhagat^{1,3}).

If ψ_1 be an L^2 -solution of (2), we must have

$$p_0^{-1}(L_{11} + m_{11} \, S_{11} + m_{12} \, S_{12}) = 0$$

$$q_0^{-1} \, (L_{21} + m_{11} \, S_{21} + m_{12} \, S_{22}) = 0$$

$$L_{11} + m_{11} \, S_{11} + m_{12} \, S_{12} = 0$$

that is,

$$L_{21} + m_{11} S_{21} + m_{12} S_{22} = 0$$

(40)

since $1/2i \exp [i\xi(x) + K_1p_1(x)q_1(x)] z(x)^{3/4} \to \infty$ as $x \to \infty$.

Similarly, if ψ_2 be an L^2 -solution of (2) we have

$$L_{12} + m_{21} S_{11} + m_{22} S_{12} = 0$$

$$L_{22} + m_{21} S_{21} + m_{22} S_{22} = 0.$$
(41)

It follows from (40-41) that

$$m_{rs}(\lambda) = \frac{U_{rs}(\lambda)}{D(\lambda)}$$
 where $U_{rs}(\lambda), D(\lambda)$

are functions of λ_1 expressible in terms of L_{rs} , $S_{rs}(\lambda)$. It follows by argument similar to Paladhi² pp. 457-458) that $U_{rs}(\lambda)$, $D(\lambda)$ are integral functions of λ . Therefore, $m_m(\lambda)$ are meromorphic functions of λ over $(-\infty, \infty)$.

The spectrum of the system (2–4) is, therefore, a countably infinite point-spectrum in the interval $(-\infty, \infty)$.

Acknowledgement

The material presented in this paper is taken from the project work completed by the first author and supported by UGC (India) Grant F.8-5(21)/87 SR(IL) under the scheme 'Minor Research Projects for Science Subject'. The authors are also grateful to the referees for their criticism and suggestions for improvement of the manuscript.

References

1. Bhagat, B .	Eigenfunction expansions associated with a pair of second-order dif- ferential equations, Ph.D. Thesis, Patna University, Patna, 1966.
2. Paladhi, B.R.	A spectral theorem for a pair of second-order differential system, J. Indian Inst. Sci., 1975, 57, 442-460.
3. Bhagat, B.	A spectral theorem for a pair of second-order differential equa- tions, Q. J. Math. (2), 1970, 21, 487-495.
4. Chakravorty, N.K.	Some problems in eigenfunction expansion (I), <i>Q. J. Math.</i> (2), 1965, 16 , 135-150.
5. Chakravorty, N.K.	Some problems in eigenfunction expansion (III), <i>Q. J. Math.</i> (2), 1968, 19 , 397-415.
6. TITCHMARSH, E.C.	Eigenfunction expansion associated with second-order differential equations, Part I, 1962, Clarendon Press.
 CONTE, S.D. AND SANGREN, W.C. 	An asymptotic solution for a pair of first-order equations, Proc. Am. Math. Soc., 1953, 4, 696-702.